

the unit cell), T-192 (tetragonal, with 190-192 atoms per unit cell), and γ -boron (high pressure phase, orthorhombic, with 28 atoms per unit cell). The new phase turned out to be a key to understanding the phase diagram of boron—the only element for which the phase diagram was unknown since its discovery 200 years ago.

Here, we report the synthesis of γ - and T-192 boron from β -boron at pressures up to 18 GPa and temperatures up to 2200 °C using a multianvil apparatus combined with x-ray diffraction (XRD) and Raman spectra. Based on the XRD and Raman results, we give the phase boundary of β -, γ -, and T-192 boron. Fig.1 shows the phase relations between β -boron (open circles), γ -boron (solid circles) and T phase (inverse triangles) based on the results of the multianvil quenched experiments. The semi-solid circles represent β -boron and γ -boron in coexistence. The line is a phase boundary between β -boron and γ -boron, and the inset show the theoretical phase boundary from Oganov *et al.*[1] and the tentative phase boundary from Zarechnaya *et al.*[2]. Additionally, the two open inverse triangle represent P-T conditions of T-192 phase from Oganov *et al.*[1] and Ma *et al.*[3] respectively. Combined with the previous results [1], [2] and our study, γ -boron phase becomes stable under a certain pressures (above \sim 8.5 GPa), and β -boron can transform into γ -boron above \sim 8.5 GPa and using heating to overcome kinetic barriers, and the kinetic barriers decrease with increasing pressure. However, at higher temperatures, β -boron and T-192 phase are more stable than γ -boron, thus γ -boron transforms back to β -boron (\sim 9 GPa) or continues to transform into T-192 phase (above \sim 10 GPa) with increasing temperature depend on undergoing high pressure.

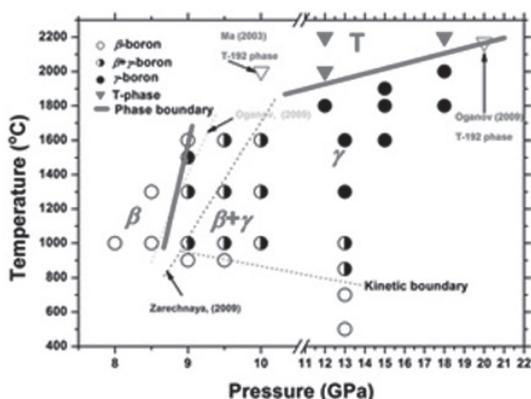


Fig.1 Phase relations between β -boron, γ -boron and T phase

[1] A.R. Oganov, J. Chen, C. Gatti, Y.Z. Ma, Y.M. Ma, C.W. Glass, Z. Liu, T. Yu, O.O. Kurakevych, V.L. Solozhenko, *Nature* **2009**, *457*, 863-867. [2] E.Y. Zarechnaya, L. Dubrovinsky, N. Dubrovinskaia, Y. Filinchuk, D. Chernyshov, V. Dmitriev, N. Miyajima, A. El Goresy, H.F. Braun, S. Van Smaalen, I. Kantor, V. Prakapenka, M. Hanfland, A.S. Mikhaylushkin, I.A. Abrikosov, S.I. Simak, *Phys. Rev. Lett.* **2009**, *102*, 185501. [3] Y.Z. Ma, C.T. Prewitt, G.T. Zou, H.K. Mao, R.J. Hemley, *Phys. Rev. B.* **2003**, *67*, 174116

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Structural anomaly in a novel iron-based perovskite

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Recent discovery of iron-based perovskites $ACu_3Fe_4O_{12}$ ($A = Ca$ and La) has been attracting much interest due to the fascinating and unexpected properties. A charge disproportionation (CD) of Fe^{4+} into Fe^{3+} and Fe^{5+} occurs in $CaCu_3Fe_4O_{12}$ (CCFO) [1], whereas an intersite charge transfer (CT) between Cu and Fe results in a large volume change in $LaCu_3Fe_4O_{12}$ [2]. The different electronic phases resulted from CD/CT imply that further anomalous properties can be achieved in this system.

A novel iron-based perovskite $SrCu_3Fe_4O_{12}$ (SCFO) was successfully synthesized using high-pressure of 15 GPa. The structural and physical properties of SCFO were in contrast to those of the known $ACu_3Fe_4O_{12}$ ($A = Ca$ and La) perovskites. SCFO demonstrated a large negative thermal expansion (NTE) with a linear expansion coefficient (ca. $-2 \times 10^{-5} K^{-1}$ at maximum) in a temperature range of 170–270 K. The Rietveld refinement based on the synchrotron X-ray powder diffraction data revealed that the NTE was attributed to a continuous intersite CT between Cu and Fe. Mössbauer spectroscopy exhibited that SCFO resulted in a charge disproportionated state below \sim 200 K. The relative abundance of $Fe^{3+} : Fe^{5+} = 4 : 1$, which is different from the ratio of 1 : 1 for CCFO, implies the electron doping into Fe through intersite charge transfer.

[1] I. Yamada *et al.*, *Angew. Chem. Int. Ed.* **2008**, *47*, 7032–7035. [2] Y.W. Long *et al.*, *Nature* **2009**, *458*, 60–63.

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Electron and magnetic properties in high temperatures magnetic semiconductors at high pressure up to 7 GPa

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In high-temperature ferromagnetic semiconductors $Cd_{1-x}Mn_xGeAs$ ($x=0\div 0.36$) and $Cd_{1-x}Mn_xGeP_2$ ($x=0\div 0.225$) there is carried out a complex investigation of electric and magnetic properties. The baric dependences of the specific resistance ρ , Hall coefficient R_H , and relative magnetic susceptibility χ/χ_0 are measured. The $\rho(P)$ and $R(P)$ are measured in high-pressure device of “Toroid” type [1], [2] when pressure rises and falls up to 7 GPa. The magnetic susceptibility is estimated by a method described in the work [3]. Structural phase transitions are found in baric dependences of $\rho(P)$ and $R_H(P)$ in both compounds at increase and decrease in pressure. A position of phase transitions sifts towards the high pressures when a percentage of Mn increases. All phase transitions are reversible in $Cd_{1-x}Mn_xGeAs_2$, in $Cd_{1-x}Mn_xGeP_2$ samples with $x \leq 0.135$ the phase transition is accompanied by partial decomposition of a substance, what confirms the X-ray diffraction study before and after pressure applying on dependences $(\chi/\chi_0)P$. In all samples of both compounds there are observed the magnetic phase transitions which shift towards high pressures with increase in percentage of Mn. When pressure decreases the hysteresis emerges. A magnetic phase transition is not revealed in base samples of $CdGeAs$ and $CdGeP$. We interpret the observed phase transitions as non-magnetic phase transition [4]. The temperature dependences of normal and abnormal Hall coefficients are calculated from magnetic-field dependences of Hall resistance for $Cd_{1-x}Mn_xGeAs_2$ ($x=0\div 0.36$) by the method of interactive graphical plotting.